TABLE III

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でので	Ec.1010, erg/g	Pc.1013, bars	PH-1011 bars
1.07 1.15 1.19 1.25 1.30 1.35 1.40 1.45 1.50 1.55 1.60 1.70	0.0525 0.181 0.30 0.585 0.86 1.18 1.57 1.97 2.42 2.94 3.51 4.12 4.79 5.54	0.132* 0.200* 0.332** 0.463 0.600 0.761 0.942 1.154 1.358 1.613 1.880 2.174 2.484 2.822	0.484 0.645 0.840 1.07 1.36 1.67 2.08 2.55 3.13 3.82 4.66
	om Ref. 8.		

sibility at high shock compression densities, where the dynamic adiabate differs considerably from the isotherm at absolute zero. Here and hereinafter, we shall not distinguish between isotherms at room temperature and at absolute zero.

We shall write the equation of state and an expression for the internal energy in the form

$$P = -\partial E_{c}/\partial v + BT/v; \qquad (10)$$

$$E = E_{\rm c} + C_{v}T. \tag{11}$$

Here $-\partial E_{C}/\partial v = P_{C}$ is the compression pressure at $T = 0^{\circ}K$, B is the thermal pressure coefficient, C_{V} is the specific heat at constant volume. In general, B and C_{V} can depend on temperature and density.

When temperature is eliminated we obtain the well-known caloric equation of state⁹⁻¹⁰

$$P + \frac{\partial E_{\mathbf{c}}}{\partial v} = \gamma \frac{E - E_{\mathbf{c}}}{v} \,. \tag{12}$$

The left member of (12) is the thermal component of the pressure and the ratio $(E - E_c)/v$ in the right member is the volumetric density of thermal energy.

The Grüneisen coefficient $\gamma = B/C_V$ gives the ratio of the thermal pressure to the thermal energy density. We shall consider this quantity to be a constant which like the unknown function $E_C(v)$ must be obtained from experiment. For this purpose, substituting in (12) the expression for E given by the Hugoniot equation (2), we obtain the differential equation.

$$v\frac{\partial E_{c}}{\partial v} + \gamma E_{c} = -\gamma \frac{P_{H}(v)v}{2} \left(1 + \frac{2}{\gamma} - \frac{v_{0}}{v}\right). \tag{13}$$

When $v_1 = v_0$, $E_c = \partial E_c / \partial v = 0$. In (13) $P_H(v)$ is the experimental equation of the shock adiabate. The solution of (13) is given by

$$E_{c} = \frac{1}{(h-1) v^{2/(h-1)}} \int_{v_{0}}^{v} P_{H}(x) \left(\frac{v_{0}}{x} - h\right) x^{2/(h-1)} dx, \quad (14)$$

$$P_{c} = -\frac{\partial E_{c}}{\partial v} = \frac{2}{(h-1)^{2}v^{(h+1)/(h-1)}}$$
(14a)

$$\times \int_{v}^{v_0} P_{\mathrm{H}}(x) \left(\frac{v_0}{x} - h \right) x^{2 \left((h-1) \right)} dx - \frac{1}{h-1} P_{\mathrm{H}}(v) \left(\frac{v_0}{v} - h \right).$$

In (14) $h = (2/\gamma) + 1$ is the so-called maximum density of shock compression.

At pressures from 3×10^5 to 5×10^6 atm, $P_{\rm H}(v)$ for iron is given by (9). We shall be committing only a very small error if, when calculating $E_{\rm C}$ in the region $P>3\times10^5$ kg/cm², we consider (9) to apply also to the initial portion of the dynamic adiabate.

In order to determine γ and thus h, we shall compare two states 1 and 2 representing the shock compression of solid and of porous iron, respectively, to the same specific volume v_1 (Fig. 1). Since both states have the same energy E_C , which depends only on the volume, the pressure difference $\Delta P = P_2 - P_1 = \Delta P_T$ is accounted for by the thermal energy difference $\Delta E = E_2 - E_1 = \Delta E_T$. From (2)

$$\Delta E_{\tau} = \frac{1}{2} \left[P_2 \left(v_{00} - v_1 \right) - P_1 \left(v_0 - \dot{v_1} \right) \right]. \tag{15}$$

Here v_{00} is the initial volume of the porous iron and P_2 is the shock compression pressure of the porous iron. The constants γ and h are obtained on the basis of (15) from the equations

$$\frac{1}{\gamma} = \frac{\Delta E_{\tau}}{v_1 \Delta P_{\tau}} = \frac{1}{2} \left[\frac{v_{00} P_2 - v_0 P_1}{(P_2 - P_1) v_1} - 1 \right]; \tag{16}$$

$$h = \frac{(v_{00} / v_0) P_2 - P_1}{P_2 - P_1} \left(\frac{v_0}{v_1}\right). \tag{16a}$$

The available experimental data on the dynamic compressibility of porous iron (in Table II) permit us to obtain γ and h for two degrees of shock compression. The values of these parameters are given in the last two columns of Table II. For the transition from the dynamic adiabate to the isotherm T=0, h is more important at high compressions, that is, in the region where the thermal components of the pressures are relatively large. In subsequent calculations, we shall assume h=2.25.

The values of the cold compression energy E_c and of the pressure $P_c = -\partial E_c/\partial v$ calculated from (14) are given in Table III for pressures from 0.3 to 5 million atmospheres. Also given are the pressures $P_H(v_0/v)$ of the shock adiabate for the compression of iron of normal density, which were calculated from the interpolation formula (9).

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